

Pressure dependence of the Shubnikov-de Haas oscillation spectrum of β'' -(BEDT-TTF)₄(NH₄)[Cr(C₂O₄)₃].DMF

David Vignolles¹, Vladimir N. Laukhin^{2,3}, Alain Audouard¹, Marc Nardone¹, Tatyana G. Prokhorova⁴, Eduard B. Yagubskii⁴ and Enric Canadell³

¹ Laboratoire National des Champs Magnétiques Pulsés ^a, 143 avenue de Rangueil, 31400 Toulouse, France

² Institució Catalana de Recerca i Estudis Avançats (ICREA), 08010 Barcelona, Spain

³ Institut de Ciència de Materials de Barcelona (ICMAB - CSIC), Campus UAB, 08193 Bellaterra, Catalunya, Spain

⁴ Institute of Problems of Chemical Physics, Russian Academy of Sciences, 142432 Chernogolovka, MD, Russia

Received: February 6, 2008/ Revised version: date

Abstract. The Shubnikov-de Haas (SdH) oscillation spectra of the β'' -(BEDT-TTF)₄(NH₄)[Cr(C₂O₄)₃].DMF organic metal have been studied in pulsed magnetic fields of up to either 36 T at ambient pressure or 50 T under hydrostatic pressures of up to 1 GPa. The ambient pressure SdH oscillation spectra can be accounted for by up to six fundamental frequencies which points to a rather complex Fermi surface (FS). A noticeable pressure-induced modification of the FS topology is evidenced since the number of frequencies observed in the spectra progressively decreases as the pressure increases. Above 0.8 GPa, only three compensated orbits are observed, as it is the case for several other isostructural salts of the same family at ambient pressure. Contrary to other organic metals, of which the FS can be regarded as a network of orbits, no frequency combinations are observed for the studied salt, likely due to high magnetic breakdown gap values or (and) high disorder level evidenced by Dingle temperatures as large as $\simeq 7$ K.

PACS. 71.18.+y Fermi surface: calculations and measurements; effective mass, g factor – 71.20.Rv Polymers and organic compounds – 72.20.My Galvanomagnetic and other magnetotransport effects

1 Introduction

The family of isostructural monoclinic charge-transfer salts β'' -(BEDT-TTF)₄(A)[M(C₂O₄)₃]·Solv have been widely studied in the past decade [1]. In the above formula, BEDT-TTF stands for bis(ethylenedithio)tetrathiafulvalene, A is a monovalent cation (A = H₃O⁺, K⁺, NH₄⁺, etc.), M is a trivalent cation (M = Cr³⁺, Fe³⁺, Ga³⁺, etc.) and Solv is a solvent molecule such as benzonitrile (C₆H₅CN), dimethylformamide (C₃H₇NO), nitrobenzene (C₆H₅NO₂) and pyridine (C₅H₅N), labelled hereafter BN, DMF, NB and P, respectively. In the following, the compounds belonging to this family are referred to as A-M·Solv. The interest in this family of compounds has been motivated by the observation of superconductivity at ambient pressure in the H₃O-Fe·BN salt ($T_c = 8.5$ K) [2]. Later on, other superconducting salts with magnetic ions were reported for this family [3,4]. Besides, a metallic and ferromagnetic ground state was achieved in the (BEDT-TTF)₃[MnCr(C₂O₄)₃] compound [5]. Contrary to the orthorhombic compounds with the same generic formula and β'' packing, which are semiconductors [6,7], all the monoclinic salts of this family exhibit metallic conductivity around room temperature. Nevertheless, a large variety of temperature-dependent behaviours and various ground states are observed which might be connected to details of their electronic structure.

According to band structure calculations [8], the Fermi surface (FS) of NH₄-Fe·DMF and (NH₄)_{0.75}K_{0.25}-Cr·DMF

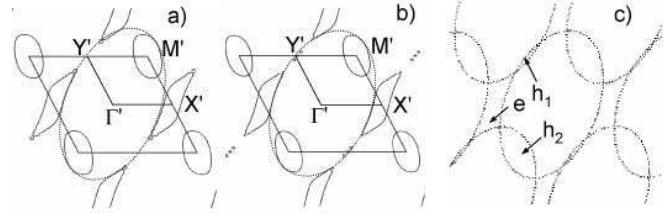


Fig. 1. Fermi surface (FS) of (a) NH₄-Fe·DMF and (b) (NH₄)_{0.75}K_{0.25}-Cr·DMF according to band structure calculations [8] in which the FS is considered on the basis of a unit cell with vectors $a' = a$ and $b' = (a + b)/2$. (c) schematic representation of intersecting elliptic hole tubes leading to three compensated electron (e) and hole (h_1 and h_2) orbits. The area of the ellipses in dotted lines (⊙ orbits, see text) is equal to that of the First Brillouin zone.

salts originates from quasi two-dimensional (2D) hole elliptic orbits, labelled ⊙ in the following, whose cross section is equal to the first Brillouin zone (FBZ) area (see Fig. 1). In the case of NH₄-Fe·DMF, these orbits intersect

^a UMR 5147: Unité Mixte de Recherche CNRS - Université

along the $(a'^* + b'^*)$ direction¹, leading to one electron and one hole compensated orbit with a cross section area of 8.8 percent of the FBZ one. Analogous FS topology has also been reported for the superconducting H₃O-Fe·BN salt [6]. These calculations are in agreement with the Shubnikov-de Haas (SdH) oscillation spectrum of the H₃O-Ga·NB salt for which only one frequency was reported [9]. Nevertheless, as pointed out in Ref. [8], the \odot orbits may also intersect in the b'^* direction leading to one or more additional orbits around the Y' point of the FBZ [see Figs. 1(b) and (c)]. This picture holds for the NH₄-Fe·DMF salt, for which the SdH oscillation spectra can be interpreted on the basis of three compensated orbits with cross section areas of 1.2, 4.8 and 6 percent of the FBZ area that are therefore connected by a linear combination [10]. However, the FS of other compounds of this family may be more complicated since four frequencies corresponding to orbit's area in the range 1.1 to 8.5 percent of the FBZ area were reported for the H₃O-M·P (M = Cr, Ga, Fe) salts [11]. In this latter case, a density wave ground state, responsible for the observed strongly non-monotonous temperature dependence of the resistance, has been invoked in order to account for this discrepancy. However, only two frequencies were observed for the H₃O-M·NB (M = Cr, Ga) salts [12]. Additional combination frequencies, typical of coupled 2D orbits networks, linked to the field-induced

chemical potential oscillation [13] and (or) field-dependent Landau level broadening [14] were also reported [10, 11].

An important feature of the oscillatory spectra of most of these compounds is the strong field-damping factor. Indeed, Dingle temperature values (T_D) in the range 2 K to 4 K were reported for e.g. H₃O-M·P [11] which is the signature of a significant disorder. This feature is in line with structural data [15, 16] which indicate that terminal ethylene groups of some of the BEDT-TTF molecules exhibit a large solvent-dependent positional disordering. As for the compounds with the DMF solvent, the DMF molecules themselves are also disordered [8]. As a matter of fact, even larger Dingle temperatures ($T_D \approx 4$ K to 6 K) were reported for the NH₄-Fe·DMF salt [10].

In order to get some insight in possible connection between the FS topology and the ground state, we report on the pressure dependence of the SdH spectrum of the NH₄-Cr·DMF salt which, contrary to the above mentioned NH₄-Fe·DMF salt exhibits a metallic conductivity down to about 10 K. Even though the overall behaviour of the resistivity as the temperature varies is unaffected by applied pressures up to 1 GPa, which suggests that the ground state remains unchanged in this pressure range, it is shown that the FS topology is very sensitive to applied pressure.

2 Experimental

The crystals studied, labelled # 1 to # 3 in the following, were elongated hexagonal platelets with approximate dimensions $(0.6 \times 0.4 \times 0.25)$ mm³ for crystal # 2 and $(0.4$

¹ In Ref. [8], the FS is considered on the basis of a unit cell with vectors $a' = a$, $b' = (a + b)/2$ and $c' = c$. This unit cell contains four BEDT-TTF molecules.

$\times 0.2 \times 0.1$) mm³ for crystals # 1 and # 3, the largest faces being parallel to the conducting *ab*-plane. Magnetoresistance experiments were performed in pulsed magnetic field of up to 36 T for crystals # 1 and # 2 and 50 T for crystal # 3, with pulse decay duration of 0.78 s and 0.32 s, respectively, in the temperature range from 2 K to 4.2 K. For crystals # 1 and # 3, the magnetic field was applied normal to the conducting plane whereas a sample holder rotating about an axis perpendicular to the magnetic field allowed a change of the direction of the magnetic field with respect to the conducting plane for crystal # 2. Crystal # 3 was studied under hydrostatic pressure in an anvil cell designed for isothermal measurements in pulsed magnetic fields [17]. The pressure applied at room temperature was estimated from the primary pressure value calibrated beforehand with a manganin piezoresistive sensor. The maximum pressure reached in the experiments was 1 GPa at low temperature, taking into account a pressure decrease of 0.1 GPa on cooling. Electrical contacts to the crystal were made using annealed platinum wires of 20 μ m in diameter glued with graphite paste. Alternating current (1 to 17 μ A, 20 kHz) was injected parallel to the *c** direction (interlayer configuration). A lock-in amplifier with a time constant of 100 μ s was used to detect the signal across the potential contacts.

Analysis of the oscillatory magnetoresistance is based on discrete Fourier transforms and direct fittings of the magnetoresistance data. Discrete Fourier transforms are calculated with a Blackman window in a given field range from B_{min} to B_{max} . The absolute value of the amplitude

(A_i) of the Fourier component with frequency F_i is determined, for a mean field value $B_{mean} = 2/(1/B_{min} + 1/B_{max})$, from the amplitude of the discrete Fourier transform (A_{calc}) as $A_i = 4A_{calc}/0.84(1/B_{min} - 1/B_{max})$. The direct fitting method was the following: a low order polynomial (typically 2th order) together with one Fourier component with initial frequency close to one of the frequencies detected in experimental data is introduced. In the case where residuals still exhibit oscillatory features, a subsequent oscillation is introduced with a frequency close to one of the main frequencies detected in the residuals. The procedure is repeated until either not any oscillatory component can be detected in the residuals or the oscillatory component corresponds to high order harmonics of already detected frequencies. Finally, the order of the polynomial, which account for the background, is increased up to at most the 4th order. It has been checked that neither the detected frequencies depend appreciably on the order in which the various components are introduced in the fittings nor the magnitude of the background significantly depends on the order of the polynomial.

3 Results and discussion

The zero-field temperature dependence of the resistance of the studied crystals is displayed in Figure 2. A metallic behaviour is observed down to about 10 K in the pressure range explored. The crystal-dependent resistance ratio $R(10 \text{ K})/R(293 \text{ K})$ is in the range $0.15 \div 0.5$. This behaviour, which is in agreement with the data of [8], is at variance with the strongly non-monotonic temperature de-

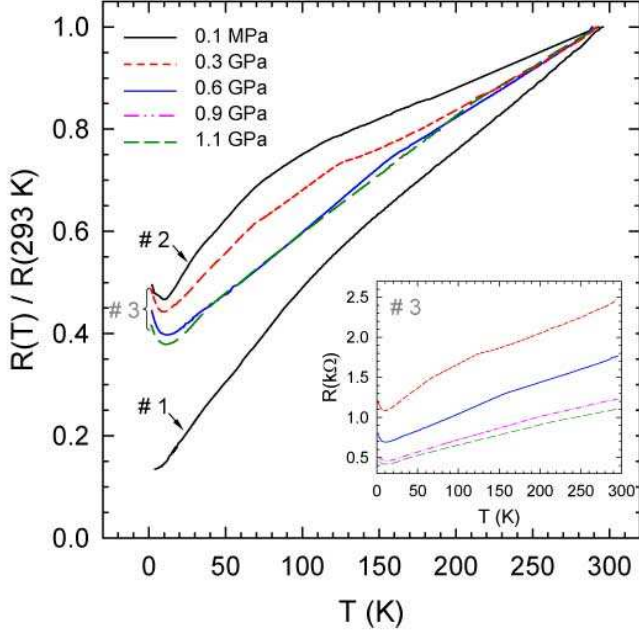


Fig. 2. Temperature dependence of the zero-field resistance normalized to the room temperature value for various pressures measured at room temperature. The crystal number (see text) is indicated on the curves. The inset displays the temperature dependence of the crystal # 3 resistance for the same pressures.

pendence reported for other salts such as $\text{H}_3\text{O-M-Solv}$ ($\text{M} = \text{Ga, Cr; Solv} = \text{NB, P}$) [9, 15], $\text{NH}_4\text{-Fe-DMF}$ [10] or $\text{H}_3\text{O-M-P}$ ($\text{M} = \text{Ga, Fe}$) [11]. In the latter case, a metal-density wave transition has been suggested to occur around 150 K in order to account for the observed behaviour. Below 10 K, a small resistance rise is observed that might be linked to disorder. Although a significant pressure dependence of the resistance is observed (see the inset of Figure 2) in

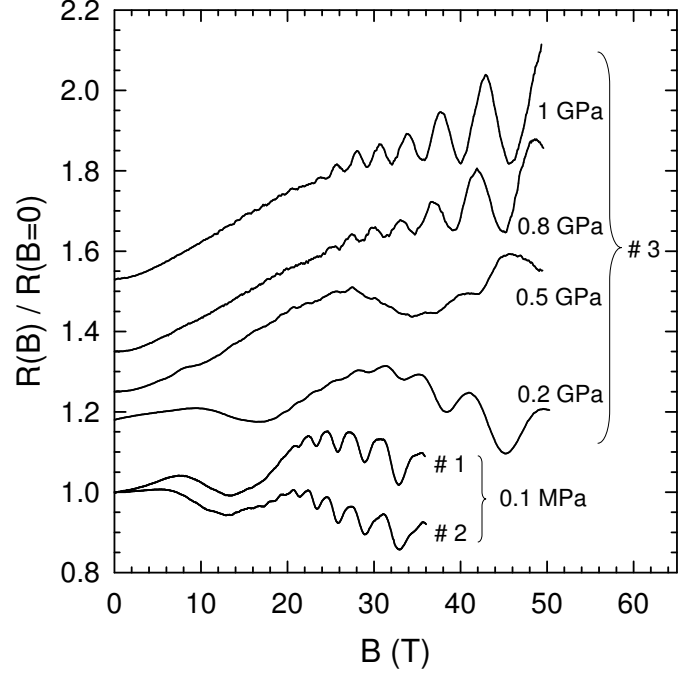


Fig. 3. Magnetoresistance at 2 K for various pressures (the low temperature values of which are given). The crystal number (see text) is indicated. Curves have been shifted from each other for clarity.

the whole temperature range explored [e.g. $\text{dln}(R)/\text{dP} \simeq 1 \text{ GPa}^{-1}$ at room temperature], the applied pressure has only a minor effect on the resistance ratio. Magnetoresistance data at 2 K are presented in Figure 3 for various pressures. In addition to magnetoresistance oscillations, a non-monotonous behaviour which appears as a slow undulation is observed up to 0.6 GPa.

In the following, we concentrate first on the ambient pressure SdH oscillation spectra (see Section 3.1). The pressure dependence of the magnetoresistance oscillation spectra is considered in Section 3.2.

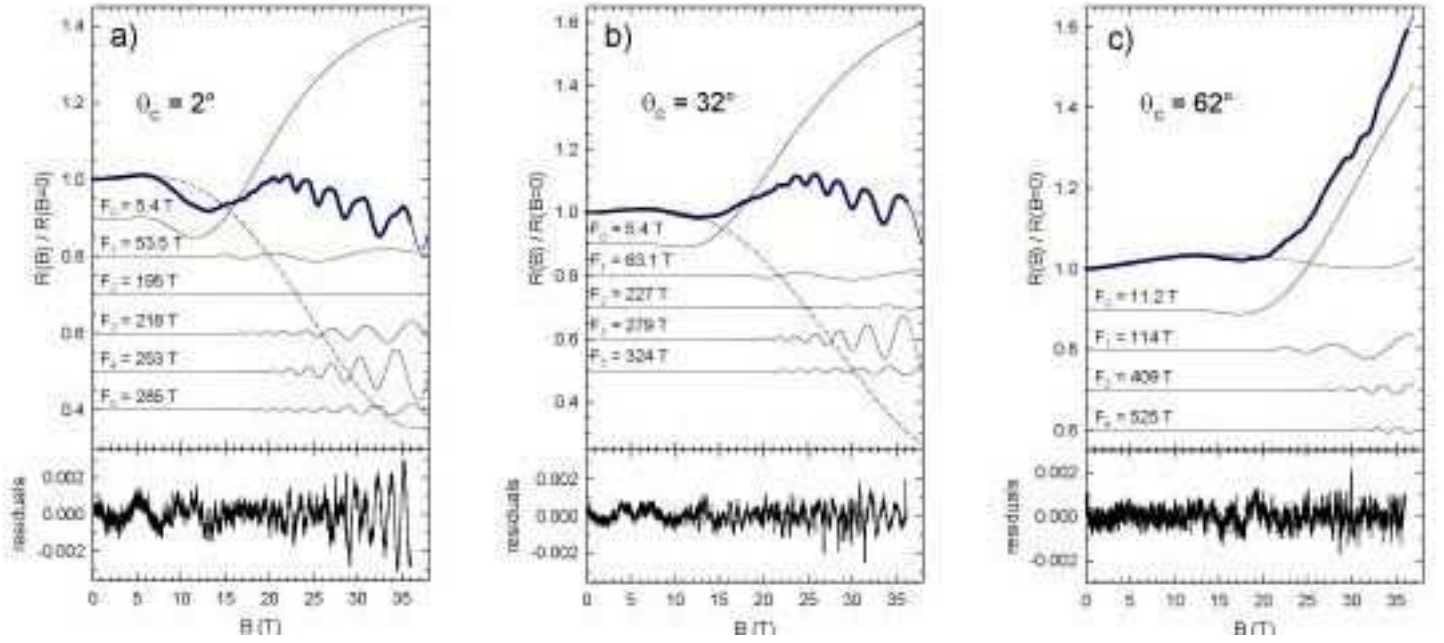


Fig. 5. Magnetoresistance data of crystal # 2 at ambient pressure (thick black solid lines). Thin blue solid lines are corresponding best fits of Eq. 6. Thin black dashed and solid lines displays the background magnetoresistance (4th order polynomial) and the contribution of each frequency (index i in Eq. 6), respectively, entering the fits. Curves have been shifted down from each other for clarity. Data are collected at 2 K for various directions of the magnetic field. Residuals values are given by $[R(B) / R(B = 0)]_{\text{experimental}} - [R(B) / R(B = 0)]_{\text{fit}}$.

3.1 Ambient pressure oscillatory spectrum

Figure 4 displays the normalized magnetoresistance of sample # 2 at a temperature of 2 K for various directions of the magnetic field. The nature of the above mentioned slow undulation, whose extremes are angle-dependent, needs clarification since it can be due to either some non-monotonic behaviour of the background magnetoresistance or to a slow SdH oscillation linked to the presence of a very small orbit. In any case, it makes the determination of the background magnetoresistance difficult. This background should be properly removed from the magnetoresistance data in order to avoid a large zero-frequency peak liable to ham-

per the extraction of reliable oscillatory data at low frequency (say below few tens of teslas) by Fourier analysis. Moreover, as developed later on and in agreement with data from other compounds of the same family [10,11], large Dingle temperatures are observed. This feature leads to a steep field dependence of the oscillation amplitude which reduces the field range in which oscillations can be detected and therefore broadens the various components' peaks appearing in the Fourier transforms. For these reasons, in addition to Fourier analysis, information on the oscillatory spectra were extracted from direct fittings of the Lifshits-Kosevich formula (LK) to the magnetoresis-

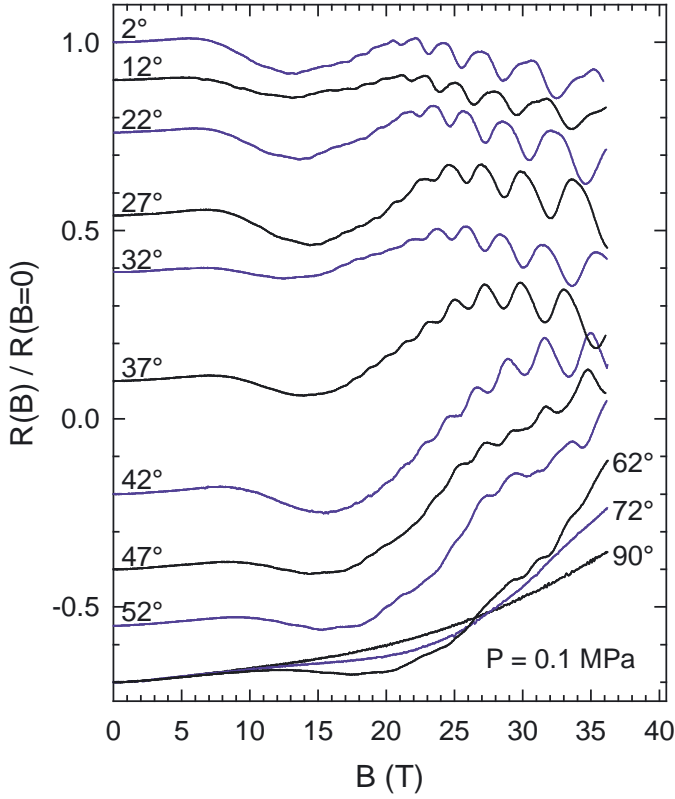


Fig. 4. Ambient pressure magnetoresistance of crystal # 2 at 2 K for various directions of the magnetic field. The angle between the magnetic field and the conducting plane is indicated. Curves have been shifted down from each other for clarity.

tance data, assuming the background magnetoresistance (R_{bg}) can be approximated by a 4th order polynomial.

According to the LK formula, the oscillatory magnetoresistance of a metal whose FS is composed of several 2D orbits is given by:

$$\frac{R(B)}{R_{bg}} = 1 + \sum_i a_i \times \sum_{\lambda=1}^{\infty} R_{Ti\lambda} R_{Di\lambda} R_{MBi\lambda} R_{Si\lambda} (-1)^{\lambda+1} \cos[2\pi\lambda(\frac{F_i}{B} - \gamma_i)] \quad (1)$$

where λ is the harmonic order. F_i and γ_i are the frequency and the phase factor of the oscillation linked to the orbit i . Equation 1 assumes that the oscillations am-

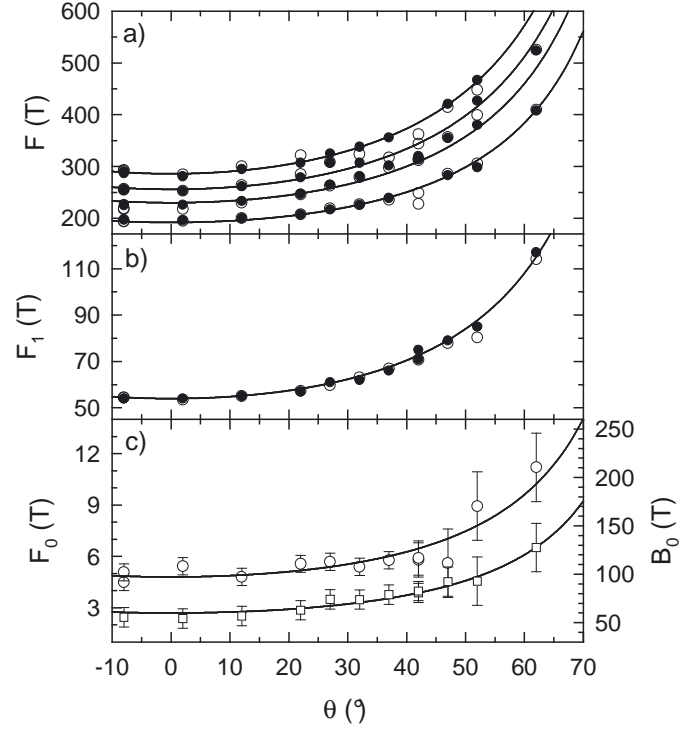


Fig. 6. Angle dependence of the frequencies observed at ambient pressure. (a) and (b) displays F_2 to F_5 and F_1 , respectively. Open and closed symbols are deduced from best fits of Eq. 6 (see Fig. 5) and Fourier analysis, respectively. (c) displays the angle variation of F_0 (circles) and of the parameter B_0 (squares) deduced from best fits of Eq. 6 to the data.

plitude is small so that $R(B)/R_{bg} - 1 \simeq 1 - \sigma(B)/\sigma_{bg}$ and that the Hall effect is either negligible (which is actually the case for interlayer magnetoresistance measurements with the magnetic field normal to the conducting plane) or contributes to the background magnetoresistance, only. The thermal (for a 2D FS), Dingle, magnetic breakthrough (MB) and spin (S) damping factors are respectively given by [18]:

$$R_{Ti\lambda} = \frac{\alpha T \lambda m_i^*(\theta = 0)}{B \cos \theta \sinh[\alpha T \lambda m_i^*(\theta = 0)/B \cos \theta]} \quad (2)$$

$$R_{Di} = \exp[-\alpha T_D \lambda m_i^*(\theta = 0)/B \cos \theta] \quad (3)$$

$$R_{MBi\lambda} = \exp\left(-\frac{t_i B_{MB}}{2B \cos\theta}\right) \left[1 - \exp\left(-\frac{B_{MB}}{B \cos\theta}\right)\right]^{b_i/2} \quad (4)$$

$$R_{Si\lambda} = |\cos(\pi\lambda\mu/\cos\theta)| \quad (5)$$

where $\alpha = 2\pi^2 m_e k_B / e\hbar$ ($\simeq 14.69$ T/K), m_i^* is the effective mass normalized to the free electron mass m_e , θ is the angle between the field direction and the conducting plane, T_D is the Dingle temperature, $\mu = g^* m_i^*(\theta = 0)/2$, g^* is the effective Landé factor and B_{MB} is the MB field. Integers t_i and b_i are respectively the number of tunnelling and Bragg reflections encountered along the path of the quasiparticle. In the high $T m_i^*/B$ range, for R_{bg} close to $R(B = 0)$ and assuming large (low) B_{MB} values and $t_i = 0$ ($b_i = 0$), in which case $R_{MBi} = 1$, Equation 1 can be approximated as:

$$R(B) \simeq R_{bg} + \sum_i \frac{A_i}{B} \times \sum_{\lambda=1}^{\infty} \exp\left(-\lambda \frac{B_i}{B}\right) \cos\left[2\pi p\left(\frac{F_i}{B} - \gamma_i\right)\right] \quad (6)$$

where A_i is a field-independent parameter, including in particular the contribution of the spin damping factor, and $B_i = \alpha(T + T_D)m_i^*/\cos\theta$. It should be kept in mind that, in addition to the above mentioned approximations, some deviations of the magnetoresistance oscillations from the LK formula are observed for 2D FS's, in particular for clean crystals at low T/B values. As a consequence, Equation 6 may not yield reliable values of the A_i and B_i parameters. Nevertheless, as reported hereafter, this equation is useful in order to identify the various components entering the Fourier spectra with a restricted number of free parameters.

Examples of best fits of Equation 6 to magnetoresistance data recorded at ambient pressure are displayed in Figure 5. According to these data, six frequencies labelled F_0 to F_5 in the following, enter the oscillatory part of the magnetoresistance for magnetic field direction not too far from the normal to the conducting plane (up to $\theta \simeq 30^\circ$), the lower frequencies F_0 to F_2 being perceptible up to $\theta \simeq 65^\circ$. The residuals displayed in the bottom part of the figures either only contain high order harmonics or do not reveal any periodic component. E.g. the high field part of the residuals of data in Figure 5(a) is dominated by the 2nd harmonic of F_3 and the 3rd harmonic of F_4 . As reported in Figure 6, the deduced frequencies follow the orbital behaviour expected for a 2D FS. Remarkably, the slow undulation is accounted for by an SdH oscillation with the frequency F_0 . In addition, a clear orbital behaviour of the parameter B_0 is observed in Figure 6(c), as it is the case for the B_i parameters relevant to the other components. A strongly negative background magnetoresistance is deduced from the fits for low θ values [see Figs. 5(a) and (b)]. Although a strongly negative magnetoresistance has already been reported above the critical field in the H₃O-M·NB superconducting salts (M = Ga, Cr) [9], it should be noticed that the main part of the magnetoresistance results from the contributions of F_0 and the background which, according to the data in Figure 5, have opposite variation above ~ 15 T. In addition, the quantum limit is reached at a few tens of teslas for this frequency. In such a case, a significant error on their amplitude cannot be excluded. The frequency values deduced from the

Table 1. Frequencies, reduced to $\theta = 0^\circ$, deduced from the best fits of Eq. 6 to magnetoresistance data at ambient pressure and from Fourier analysis (see Fig. 7). i is the frequency index appearing in Eqs. 1 and 6.

i	crystal # 1		crystal # 2	
	fit	FT	fit	FT
0	6.2 ± 1.0		5.0 ± 0.5	
1	54 ± 1	54.8 ± 0.5	53.5 ± 1.5	53.7 ± 1.0
2	202 ± 3	200 ± 2	193 ± 3	196 ± 3
3	229 ± 6	230 ± 5	218 ± 7	226 ± 4
4	257 ± 5	250 ± 2	253 ± 5	253 ± 4
5	295 ± 3	287 ± 5	288 ± 6	284 ± 6

fits of the magnetoresistance data for crystals # 1 and # 2 are displayed in Table 1. Since the frequency F_0 is very low, Fourier analyses have been performed subtracting the contributions of both the background magnetoresistance R_{bg} and the oscillation with frequency F_0 from the magnetoresistance data. An example is given in Figure 7: a good agreement between Fourier analysis of the fits of Equation 6 and of the experimental data is observed.

F_0 may correspond to an orbit with a very small cross section amounting to 0.1 percent of the FBZ area only. Such a low value is compatible with band structure calculations since the cross section values of the orbits close to the Y' point of the FBZ in Figure 1(b) are of the same order of magnitude or even smaller. The values of the other observed frequencies (F_1 to F_5) correspond to orbital areas in the range 1 to 7 percent of the FBZ area. Such values are of the same order of magnitude as those deduced from the data of Refs. [10,11]. For example, F_1 is

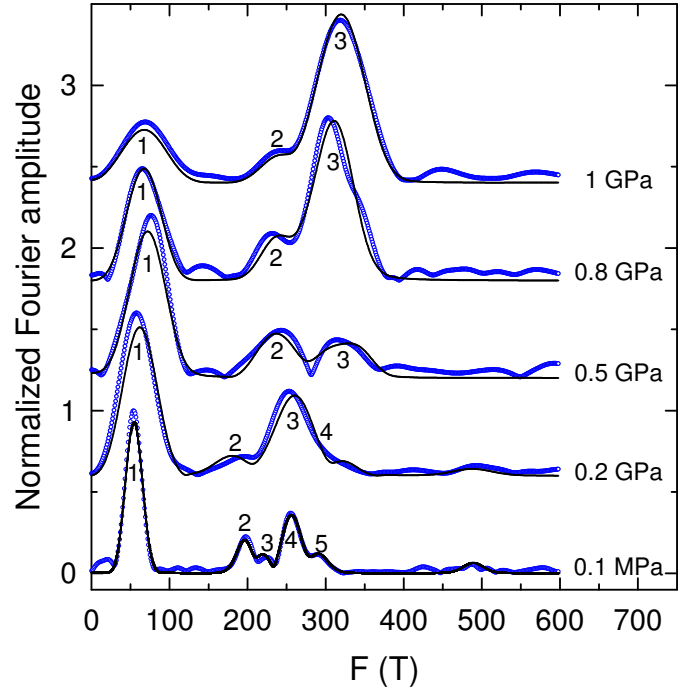


Fig. 7. Fourier spectra of the oscillatory magnetoresistance of crystals # 2 (at ambient pressure) and # 3 (under applied pressure) at 2 K (blue symbols) and of the corresponding fits of Eq. 6 (solid black lines). The magnetic field range is 8 T - 36 T for the data at 0.1 MPa, 15 T - 50 T for the data in the range 0.2 GPa to 0.8 GPa and 18 T - 50 T for the data at 1 GPa. The spectra, which are shifted from each other for clarity, have been normalized to the component with the highest amplitude. The labels correspond to the frequency index used in the text.

very close to the frequencies $F_a = 48$ T and $F_\alpha = 38$ to 50 T reported for the $\text{NH}_4\text{-Fe-DMF}$ [10] and $\text{H}_3\text{O-M-P}$ [11] salts, respectively, while F_4 is very close to the frequency $F_b = 248$ T reported for $\text{NH}_4\text{-Fe-DMF}$ [10]. It can be deduced from the data in Table 1 that $F_0 + F_1 + F_2 + F_4$ is equal to $F_3 + F_5$ within the error bars, as expected for a compensated metal. However, even in the case where this latter relationship is not fortuitous, the large number of observed frequencies cannot be fully understood on the ba-

sis of the band structure calculations displayed in Figure 1. Nevertheless, a FS based on the intersection of elliptic 2D tubes scheme might still account for a large number of orbits. If true, the actual picture would be less naive than that displayed in Figure 1(c) and, in any case, this point needs a more detailed determination of the FS topology. Other relationships such as $F_4 = F_1 + F_2$ or $F_5 = F_1 + F_4$ are also observed (see Table 1). Still in the case where they are not fortuitous, such linear combinations could indicate that frequencies F_4 and (or) F_5 are linked to either MB orbits or frequency combinations. However, since band structure calculations cannot yield detailed FS topology, no reliable conclusion can be drawn regarding the presence of MB orbits. Effective masses deduced from the temperature dependence of the Fourier component's amplitudes are in the range from 0.4 to 1.2 free electron mass (see Table 2) which is of the same order of magnitude as for other salts of this family [10,11]. Dingle temperature values are high ($T_D \sim 7$ K) which certainly rules out frequency combinations due to an oscillation of the chemical potential [13]. Otherwise, although crystals # 1 and # 2 exhibit significantly different residual resistance ratios, their Dingle temperature are rather close. This feature is in line with the statement of Ref. [11] that the crystals are composed of a mixture of insulating and metallic domains, although a metallic conductivity is observed in the present case. Within this picture, the temperature dependence of the resistance reflects the relative parts of metallic and insulating domains while the oscillatory behaviour is only related to the metallic parts which are in turn character-

ized by a large disorder as indicated by the large measured Dingle temperatures.

3.2 Pressure-dependent oscillatory spectra

Examples of best fits of Equation 6 to the data collected under applied pressure are displayed in Figure 8. As it is the case for the ambient pressure data, a good agreement with Fourier analysis is obtained (see Fig. 7). The salient feature of the pressure dependence of the oscillatory spectra is the progressive decrease of the number of observed frequencies as the pressure increases. Indeed, only five and four frequencies can be detected at 0.2 GPa and 0.5 GPa, respectively. In addition, the slow undulation attributed to F_0 cannot be detected above 0.5 GPa. Finally, only three frequencies labelled F_1 to F_3 in Figures 8(b) and (c) are observed at 0.8 GPa and 1 GPa. Both the effective masses (see Table 2) and the Dingle temperatures (see Table 2 and Fig. 10) remain constant within error bars. The three frequencies observed at high pressure are connected by the relation $F_1 + F_2 = F_3$. This point is in agreement with e.g. Fourier analysis at 1 GPa that yields $F_1 = (68 \pm 2)$ T, $F_2 = (238 \pm 4)$ T and $F_3 = (313 \pm 7)$ T, respectively. This suggests that the corresponding orbits are compensated. As it is the case for the NH₄-Fe-DMF salt [10], the high pressure spectra of NH₄-Cr-DMF can be accounted for by the band structure calculations assuming the \odot orbit, from which originates the FS, intersects both along the $(a'^* + b'^*)$ and a'^* directions yielding 3 compensated electron and hole orbits as depicted in Figure 1(c). Within this framework, the frequencies F_1 , F_2 and F_3 observed at

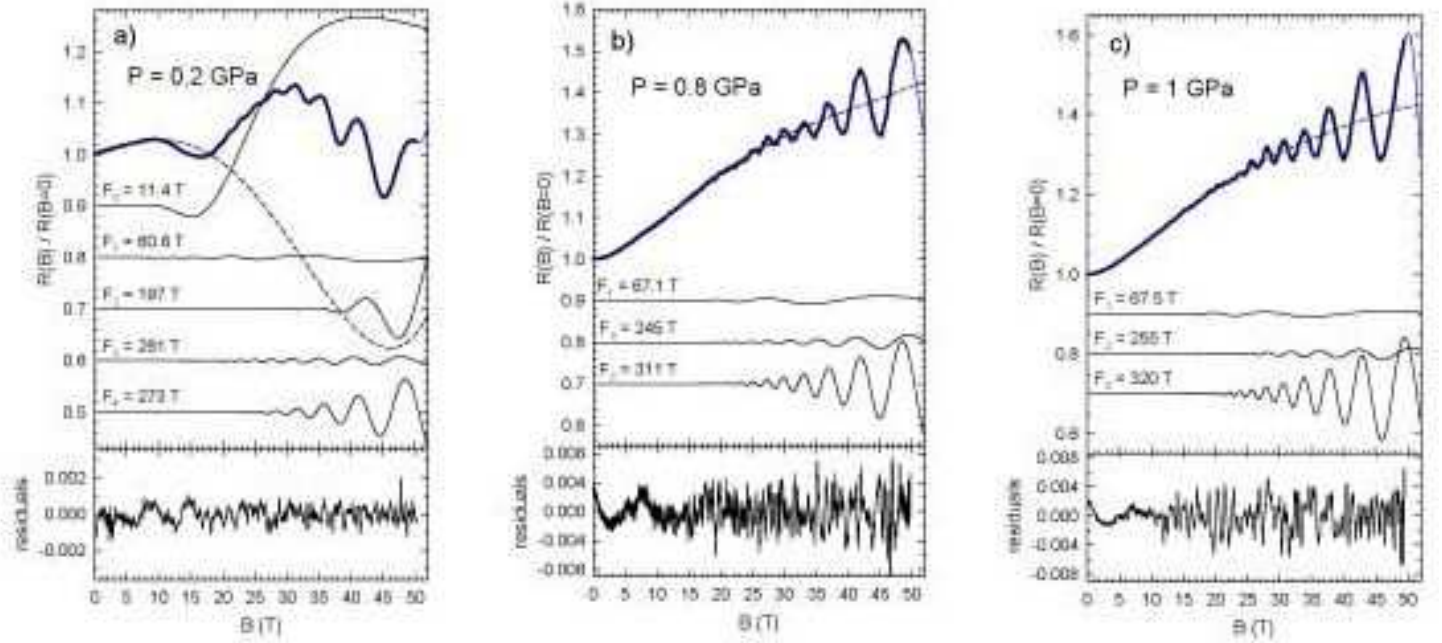


Fig. 8. Same as Fig. 5 for crystal # 3 at various applied pressures.

Table 2. Effective masses and Dingle temperatures deduced from temperature and field dependence, respectively, of the amplitude of the various oscillations observed. i is the frequency index appearing in Eqs. 1 and 6. The crystal numbers are indicated.

	0.1 MPa (# 1)		0.1 MPa (# 2)		0.5 GPa (# 3)		0.8 GPa (# 3)		1 GPa (# 3)	
i	m_i^*	T_D	m_i^*	T_D	m_i^*	T_D	m_i^*	T_D	m_i^*	T_D
1	0.4 ± 0.2	7.5 ± 5.0	0.4 ± 0.2	6.5 ± 4.0	0.5 ± 0.05	8 ± 2	0.4 ± 0.2	5.0 ± 2.5	0.60 ± 0.15	4 ± 2
2	0.4 ± 0.1	6 ± 4	0.4 ± 0.1	7 ± 4.5	0.6 ± 0.2	6.5 ± 3.5	0.80 ± 0.25	7.5 ± 4.0	0.8 ± 0.2	
3	0.7 ± 0.2		0.6 ± 0.2		0.9 ± 0.2		1.10 ± 0.25	8.5 ± 3.5	0.8 ± 0.2	9 ± 3
4	1.2 ± 0.2	6 ± 3	1.2 ± 0.2	6 ± 2						
5	0.65 ± 0.20		0.65 ± 0.20							

0.8 GPa and 1 GPa can be ascribed to the frequencies labelled F_a , F_{b-a} and F_b observed at ambient pressure in the NH₄-Fe-DMF salt.

The various frequencies observed as the pressure varies are collected in Figure 9. Provided a given orbit keeps its identity as the applied pressure varies, the pressure de-

pendence of the relevant oscillation frequency can be accounted for by a relationship of the form $d[\ln(F)]/dP = \kappa$ where κ is related to the compressibility tensor. The above relationship holds for various organic metals based on the BEDT-TTF molecule with κ ranging from 0.14 GPa⁻¹ [19] to 0.7 GPa⁻¹ [20]. The solid line in Figure

9(a) has been derived assuming this is the case for F_1 . It is obtained with $\kappa = 0.6 \text{ GPa}^{-1}$ which is within the above range. A clear downward departure from this line is nevertheless observed above 0.5 GPa. A maximum in the pressure dependence of the frequency linked to the closed orbit was also reported at $\sim 0.6 \text{ GPa}$ for β'' -(BEDT-TTF) $_2$ SF $_5$ CH $_2$ CF $_2$ SO $_3$ [21]. In this latter compound a pressure-induced phase transition is observed, although at 1.2 GPa. Regarding the other frequencies, no clear pressure dependence can be derived from data in figure 9(b) which certainly accounts for the observed drastic change of the FS topology as the applied pressure varies.

With regards to frequency combinations linked to either field-dependent Landau level broadening or chemical potential oscillations, no such features are observed, at least in the high pressure range. As a matter of fact, the frequency labelled F_{a+b} , observed in the NH $_4$ -Fe-DMF salt, which would correspond to the frequency $F_1 + F_3$ in the present case, is not observed. This can be considered at the light of the large Dingle temperatures observed ($T_D \sim 4$ to 9 K under applied pressure), keeping in mind that oscillations of the chemical potential are strongly damped by disorder. Since the other source of frequency combination is the Landau level broadening induced by coherent magnetic breakdown [14], the absence of such frequencies suggests that the FS only contains individual orbits, at least at high pressure.

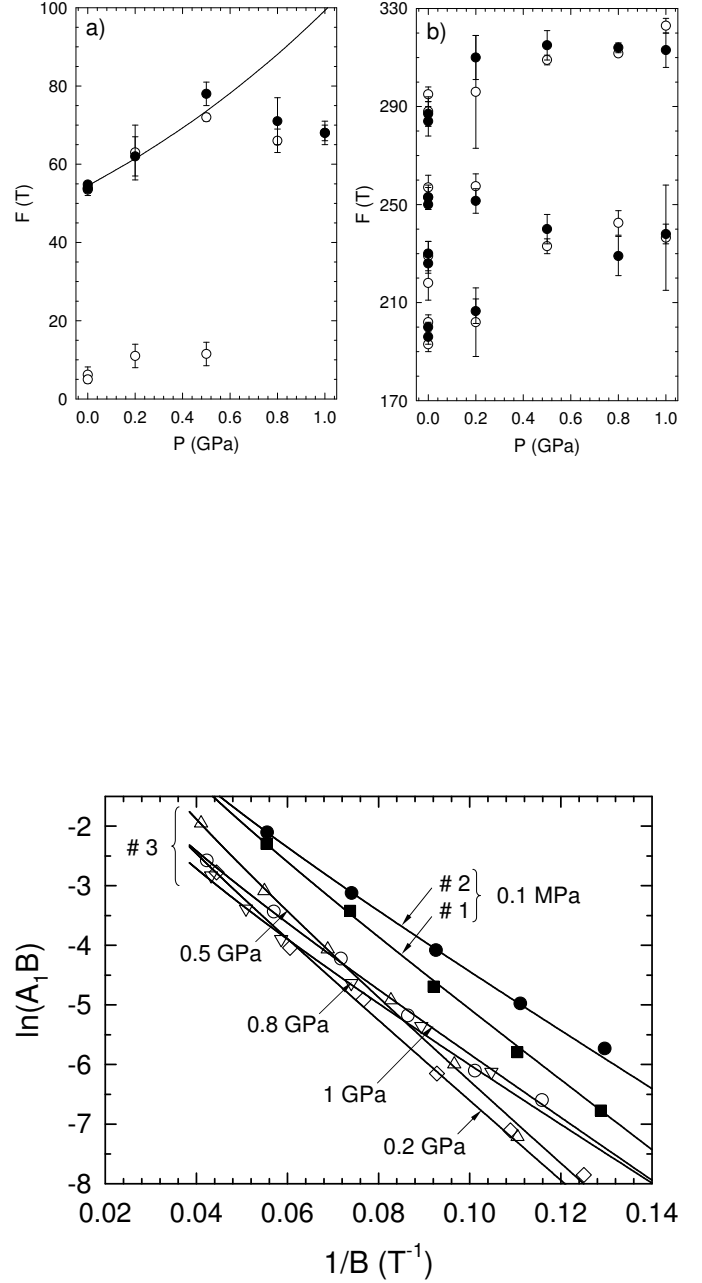


Fig. 10. Dingle plots of the Fourier amplitude A_1 at 2 K corresponding to the component at the F_1 frequency. Solid lines are best fits of Eq. 1 to the data.

4 Summary and conclusion

The pressure dependence of the FS topology of the β'' -(BEDT-TTF) $_4$ (NH $_4$)[Cr(C $_2$ O $_4$) $_3$]-DMF organic metal have been studied up to 1 GPa. The SdH oscillation spectra

observed at ambient pressure are compatible with a FS composed of up to six individual orbits. A drastic change of the FS topology is observed under pressure. As a matter of fact, the number of orbits decreases as the applied pressure increases. At 0.8 GPa and above, only three compensated orbits are observed, as it is the case for the NH₄-Fe·DMF salt at ambient pressure. This feature suggests similar FS in both cases, although the latter compound exhibits a strongly non-monotonous temperature-dependent behaviour. This result demonstrates that such non-monotonous behaviour is not necessarily connected with a density wave condensation, which was invoked in the case of H₃O-M·P [11]. At variance with magnetoresistance data of the NH₄-Fe·DMF compound and more generally of many networks of coupled orbits, no frequency combinations, due to e. g. field-induced chemical potential oscillations, were observed. This is likely connected to the large amount of disorder present in the studied crystals as indicated by the very large measured Dingle temperatures ($T_D \sim 7$ K).

This work was supported by the French-Spanish exchange programme between CNRS and CSIC (number 16 210) and by Euromagnet under the European Union contract R113-CT-2004-506239. We acknowledge Geert Rikken for interesting discussions.

References

1. E. Coronado and P. Day, *Chem. Rev.* **104** (2004) 5419
2. A. W. Graham, M. Kurmoo and P. Day, *J. Chem. Soc., Chem. Commun.* (1995) 2061.
3. L. Martin, S. S. Turner, P. Day, F. E. Mabbs and E. J. L. McInnes, *Chem. Commun.* (1997) 1367.
4. S. Rashid, S. S. Turner, P. Day, J. A. K. Howard, P. Guionneau, E. J. L. McInnes, F. E. Mabbs, J. H. Clark, S. Firth and T. Briggs, *J. Mater. Chem.* **11** (2001) 2095.
5. E. Coronado, J. R. Galán-Mascarós, C. J. Gómez-García and V. Laukhin, *Nature* **408** (2000) 447.
6. M. Kurmoo, A. W. Graham, P. Day, S. J. Coles, M. B. Hursthouse, J. L. Caulfield, J. Singleton, F. L. Pratt, W. Hayes, L. Ducasse and P. Guionneau, *J. Am. Chem. Soc.* **117** (1995) 12209.
7. L. Martin, S. S. Turner, P. Day, P. Guionneau, J. A. K. Howard, D. E. Hibbs, M. E. Light, M. B. Hursthouse, M. Uruichi and K. Yakushi, *Inorg. Chem.* **40** 1363 (2001).
8. T. G. Prokhorova, S. S. Khasanov, L. V. Zorina, L. I. Buravov, V. A. Tkacheva, A. A. Baskakov, R. B. Morgunov, M. Gener, E. Canadell, R. P. Shibaeva and E. B. Yagubskii, *Adv. Funct. Mater.* **13** (2003) 403.
9. A. Bangura, A. Coldea, A. Ardavan, J. Singleton, A. Akutsu-Sato, H. Akutsu and P. Day, *J. Phys. IV France* **114** (2004) 285.
10. A. Audouard, V. N. Laukhin, L. Brossard, T. G. Prokhorova, E. B. Yagubskii and E. Canadell, *Phys. Rev. B* **69** (2004) 144523.
11. A. Coldea, A. Bangura, J. Singleton, A. Ardavan, A. Akutsu-Sato, H. Akutsu, S. S. Turner and P. Day, *Phys. Rev. B* **69** (2004) 085112.
12. A. Bangura, A. Coldea, J. Singleton, A. Ardavan, A. Akutsu-Sato, H. Akutsu, S. S. Turner and P. Day, *Phys. Rev. B* **72** (2005) 14543.

13. A. S. Alexandrov and A. M. Bratkovsky, Phys. Rev. Lett. **76** (1996) 1308; Phys. Lett. A **234** 53 (1997) and Phys. Rev. B **63** 033105 (2001); T. Champel, ibid. **65** 153403 (2002); K. Kishigi and Y. Hasegawa, ibid. **65** 205405 (2002); J. Y. Fortin, E. Perez and A. Audouard, ibid. **71** (2005) 155101.
14. A. B. Pippard, Proc. Roy. Soc. (London) **A270** 1 (1962); P. S. Sandhu, J. H. Kim and J. S. Brooks, Phys. Rev. B **56** (1997) 11566; J. Y. Fortin and T. Ziman, Phys. Rev. Lett. **80** (1998) 3117; V. M. Gvozdkov, Yu V. Pershin, E. Steep, A. G. M. Jansen and P. Wyder, Phys. Rev. B **65** (2002) 165102.
15. H. Akutsu, A Akutsu-Sato, S. S. Turner, D. Le Pevelen, P. Day, V. N. Laukhin, A.-K. Klehe, J. Singleton, D. A. Tocher, M. R. Probert and J. A. K. Howard, J. Am. Chem. Soc. **124** (2002) 12430.
16. S. S. Turner, P. Day, K. M. Abdul Malik, M. Hursthouse, S. Teat, E. MacLean, L. Martin and S. French, Inorg. Chem. **38** 3543 (1999).
17. M. Nardone, A. Audouard, D. Vignolles and L. Brossard, Cryogenics **41** (2001) 175.
18. D. Shoenberg, Magnetic Oscillations in Metals (Cambridge University Press, Cambridge, 1984)
19. J. S. Brooks, X. Chen, S. J. Klepper, S. Valfells, G. J. Athas, Y. Tanaka, T. Kinoshita, N. Kinoshita, M. Tokumoto, H. Anzai and C. C. Agosta, Phys. Rev. B **52** (1995) 14457.
20. A. Audouard, P. Auban-Senzier, V. N. Laukhin, L. Brossard, D. Jérôme and N. D. Kushch, Europhys. Lett. **34** (1996) 599.
21. J. Hagel, J. Wosnitza, C. Pfeleiderer, J. A. Schlueter, J. Mohtasham and G. L. Gard, Phys. Rev. B **68** (2003) 104504.